### Section 2-2 - Electric Potential

[Potential](#Potential)   
[Equipotential Surfaces](#Equipotential)   
[Relationship between V and **E**](#Relationship_between_V_and)   
[Correlation to your Textbook](file:///C:/My%20Documents/finished%20web%20pages/phys543%205_30_02/apcorrelation.htm#2-2)

#### Potential

In much the same way as we used several ‘pictures’ to study motion (force and acceleration, work and energy), we will do the same this semester. We shall define the *electric potential* (V) at a given point in space to be the *electrical potential energy* (EPE) per unit of charge of a test charge located at that point; in a manner similar to that for the definition of the electric field, we are able to let the test charge diminish in size until it vanishes, leaving the quotient of our definition intact:   
V = limqTEST\_>0 EPE/qTEST.   
Note that there need not actually be a test charge at the point in order for there to exist a potential at that point, in the same way that there need not be a test charge at some point in order for an electric field to exist.

Let's do a simple example which will give us a useful result and perhaps provide a more concrete picture of the electric potential, V:   
  
Consider a point charge +Q fixed in space.  Place a small test charge +q a distance rA from Q.  There will be a repulsive force FE acting on it of magnitude keQq/r2.  Now, suppose that for whatever reason the charge q moves closer to Q, to a distance rB.  As the electric force acts though a distance s, work is done.  How much?  Can we simply multiply FE s cosF, s?

Top of Form



Bottom of Form

Note that we have found the work done by the electric force only, and have ignored any work done by other forces.

Later, we will show that the electric force is a [conservative force](#conservative); in that case, we can re-write the work done by the electric force as a change in the *electrical potential energy* of charge q:   
EPE = - WE = + keQq[1/rB - 1/rA].   
Now, we previously defined V as EPE/q, so it's safe to say that   
V = (EPE)/q = keQ[1/rB - 1/rA],   
where by V we mean VB - VA.   
So, where are we now?   
VB - VA = keQ/rB - keQ/rA.   
Think back to when you were working with gravitational potential energy (GPE).  The derivation we did then was that the change in GPE for an object moving vertically was:   
GPEf - GPEi = mghf - mghi.   
What we did next was to assign a value of zero to the GPE at some particular level, usually ground level where h = 0 as well.  In that way, we simplified this relationship to obtain   
GPEf = mghf,   
or more generally,   
GPE = mgh.   
We're about to do the same thing here.  Assign a value of zero potential to locations infinitely far from Q (that is, let VA = 0 at rA = infinity), so that the relation becomes:   
VB - 0 = Qke/rB - 0,   
or more generally   
V(r) = Qke/r.

This gives us an alternate notion of the definition of the potential: the potential at a point P is the work necessary per unit charge to bring that charge very slowly from infinity to P. The ‘very slowly’ is so that we need not consider any kinetic energy.

Note two things: the potential does not depend on the existence of the test charge q, and that this relationship is only valid for a point charge (or for regions outside of spherically symmetric charges which produce the same configuration of electric field), although similar relations can be valid for other shapes of charge.

The unit for electric potential is the *volt*.  A potential of one volt at some location means that each coulomb of charge placed there will possess one joule of EPE.

Since potential is related to work and energy, it is a scalar quantity.  The total potential at any given spot is the algebraic sum of the potentials at that spot

Vtotal = Wtotal/q = Ftotal dr/q = (Fi)d/q = ( Fi d)/q = ( Wi)/q =  (Wi/q) =  Vi   
where we realize that the work calculations may require taking dot products and doing integrals, so this derivation was very conceptual in nature.

So, if we want to find the potential caused by a collection of charges, or by a continuous distribution of charge, in principle, we need only find the potential due to each tiny bit of charge and add.

Also, note that the potential caused by a charge at some point depends on the sign of the charge.

Examples:   
Consider a ring of radius R with a charge Qo distributed evenly along its circumference.  Find the potential at the centre of the ring.

Top of Form



Bottom of Form

Consider one small charge qi, so small it looks like a point charge.  We already worked out that the potential due to a point charge is keq/r, so the contribution to the total potential from one of these little charges is   
keqi/ri.   
The total value of the potential at the centre will be the sum of all the contributions from each of the individual charges:   
V = i keqi/ri.   
In this case, all of the ris equal R, so   
V = i keqi/R = (ke/R)iqi = (ke/R) Qo = keQo/R.

What is the potential a distance x out along the axis of the ring, through the centre and perpendicular to the plane of the ring?

Top of Form



Bottom of Form

The calculation is similar to the one above, except that each charge is now a distance ri = [R2 + x2]1/2 from the point:   
V = i keqi/ri = i keqi/[R2 + x2]1/2 = (ke/[R2 + x2]1/2)iqi = (ke/[R2 + x2]1/2) Qo = keQo/[R2 + x2]1/2.

#### Equipotential Surfaces

An equipotential surface is one along which the electric potential is constant, or if you like, along which we need do no work to move a test charge.   
Consider again a positive test charge in an electric field.  The field will exert a force on q in the direction of **E**.  If the test charge moves in the direction of **E**, work is done, and the potential energy of q changes, and since V = EPE/q, the potential changes as well.  A similar argument holds if q moves against the field.  The only way that the potential will not change is if no work is be done on q, which means it must move perpendicularly to the field, *i.e.*, make the displacement perpendicular to the electric force (remember that W = Fd cosF,d).

So, we conclude that equipotential surfaces are always perpendicular to the electric field.

Now, since we know that the electric field must always be perpendicular to the surface of a metal, what can we say about the potential on the surface of a piece of metal?

Top of Form



Bottom of Form

What can you say about the potential inside of a metal?

Top of Form



Bottom of Form

#### Relationship between V and E.

There is a relationship between the potential field and the electric field which goes beyond just the direction of **E**; we can get a sense of the magnitude of **E** as well from V.  Let's once again consider a very special case which will give some insight:   
Suppose that we put a positive test charge q in a uniform electric field, **E**.   
  
There will be a force from the field on q with magnitude   
FE = qE   
acting to the right.  Let q move toward the right a distance s.  The work done by the electric field will then be   
WE = FE s = qE s.   
We know that the work done can be written as a change in EPE:   
WE = - EPE = - qV.   
When we compare these results, we see that   
qE s =  - qV   
E s = - V   
E  = (-)V/s.   
This tells us that the magnitude of **E** can be obtained by looking at how quickly the potential changes with distance along a field line.  The negative sign indicates that the **E**-field points in the direction in which V is decreasing.  Now, this result is only exactly true for the situation from whch is was developed, but conceptually, it is O.K. in almost any circumstance.  For example, if we had moved q at an angle to **E**, there would have had to have been a cosE,s term included as well: V = (-)E s cosE,s.  Also, this gives us a different set of units for **E**: V/m as well as N/C.

Here, we'll finally make the argument for **E** being a conservative field.  Consider a point charge Q.  The electric field, and therefor the electric force on a test charge q, goes as 1/r2 and is radial (inward or outward, depending on the relative signs of the charges).   Take the test charge q from Point A to Point B along Path One:   
  
Along Path 1, no work is done along the curved part, since the electric force is perpendicular to the displacemnt at all points, and there is some work done along the radial part of the path (which we won't actually calculate).  Path 2 (in red) is any other path from Point A to Point B; this path can be approximated to an arbitrary degree of accuracy with a number of radial and circular sections, as shown.  Along each circular section, no work is done.  Along each radial section, the work must be calculated, but it should be clear that work done during any outward motion from r1 to r2 will be cancelled by the work done moving radially from r2 to r1.   Only the net motion from rA to rB will result in a net amount of work done.  Since the electric force has the same magnitude and relative orientation to the displacement at all points a distance r from Q, the amount of work done will be independent of the path taken.  This last is one way of saying that the force is conservative.   
Now, what if the charge distribution had been more complicated than a point charge?  We can approximate any shape charge distribution by a collection of point charges; the work done by the E-field while the test charge moves is the sum of the works done by the individual contributions to the field of the individual charges and so it is also conservative.

[Continue to Next Section](file:///C:/My%20Documents/finished%20web%20pages/phys543%205_30_02/ap2notes3.htm)   
[Return to Directory](file:///C:/My%20Documents/finished%20web%20pages/phys543%205_30_02/apnotesdir.htm)

D Baum - 2001   
 